

# Uniaxial-stress-control of competing inter-chain exchange interactions of isosceles-triangular lattice Ising magnet CoNb<sub>2</sub>O<sub>6</sub>

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The isosceles triangular lattice Ising antiferromagnet is characterized by the ratio of exchange interactions defined as  $\gamma = J_1$  (along the base direction) /  $J_2$  (along the equilateral direction), and its magnetic property dramatically changes, depending on whether  $\gamma$  is larger than 1.0 or not. As one of the model materials, we have studied an Ising magnet CoNb<sub>2</sub>O<sub>6</sub>, where the quasi-1D ferromagnetic zigzag chains along the  $c$  axis form a frustrated antiferromagnetic isosceles-triangular lattice (ITL) with  $\gamma \simeq 1.33$  in the  $a$ - $b$  plane. If the exchange ratio  $\gamma$  can be controlled in CoNb<sub>2</sub>O<sub>6</sub> via anisotropic deformation of ITL by uniaxial pressure, variety of interesting magnetic features intrinsic to  $\gamma$  would be observed.

Actually along this context, we succeeded in crossing the Wannier point ( $\gamma=1$ ) and providing access to the region of  $\gamma < 1$  by further applying uniaxial pressure  $p$  up to 1GPa along the  $c$  axis, where AF-II magnetic ordering with  $q=1/2$  is switched to AF-I magnetic ordering with  $q=0$  at the critical pressure  $p \simeq 0.8$ GPa at which the exchange ratio  $\gamma$  becomes 1.0. This is what exactly suggested by Stephenson's exact calculation for 2D ITL, as is in the experimental reports 22 (N0 1802).

In present experiment, as a continuation of the proposal, we obtained entire  $H_{//c} - T$  magnetic phase diagram (with  $\gamma \sim 0.91$ ) at  $p = 1.1$ GPa, including determination of magnetic structure of magnetic state with  $q=1/5$  newly appearing between AF-I magnetic ground state and field induced ferrimagnetic (FR) phase, as shown in Fig.1. How the IC phase appears in  $H_{//c} - T$  magnetic phase diagram at  $p = 1.1$ GPa is entirely different from that at  $p = 0$  Pa, corresponding to that one exchange  $J_1$  (along the base direction) wins

two exchange  $J_2$  (along the equilateral direction) and there is essentially no competition among spins on isosceles triangular lattice for  $\gamma < 1$ . Level-crossing considerations using experimentally obtained four magnetic transition fields between AF-I, 5SL, FR, HHFR and IF phases as well as  $\gamma \sim 0.91$  suggest that ten times smaller negative exchange  $J_4$  and  $J'_3$  (compared with  $J_1$  and  $J_2$ ) is necessary for the appearance of 5SL and HHFR phases, respectively.

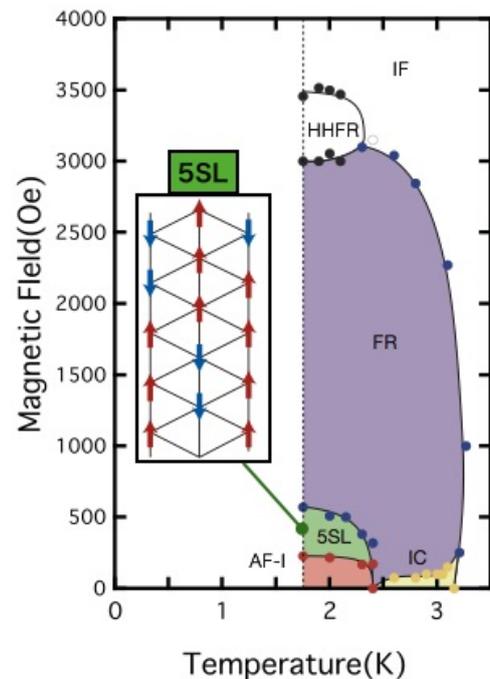


Fig. 1.  $H_{//c}$ - $T$  magnetic phase diagram at  $p = 1.1$ GPa.